

Stochastic Schrödinger equations with unbounded coefficients ¹

Carlos M. Mora

Departamento de Ingeniería Matemática - CI²MA
Universidad de Concepción

10th International Conference on Operations Research
March 6 - 9 (2012), Habana

¹Supported in part by FONDECYT Grant 1110787 and by BASAL Grants PFB-03 and FBO-16



Outline

- 1 Stochastic Schrödinger equations
 - Examples
 - General Stochastic Schrödinger equations
- 2 Basic properties
 - Relation between linear and non-linear SSEs
 - Quantum master equations
 - Well-posed of the linear SSE
- 3 OQS in position representation
 - General model
 - Ehrenfest's theorem
- 4 Regular invariant states



Example 1

Paul trap - Fluctuations in the location of the center

State space: $\mathfrak{h} = L^2(\mathbb{R}, \mathbb{C})$

Hamiltonian: $H = -\frac{1}{2M} \frac{d^2}{dx^2} + \frac{1}{2} M \omega^2 x^2$, with $M > 0$ and $\omega \in \mathbb{R}$

$L_1 = -i\eta x$, with $\eta > 0$.

$$X_t = X_0 + \int_0^t \left(-iH + \frac{1}{2} L_1^* L_1 \right) X_s ds + \int_0^t L_1 X_s dW_s^k$$

M.E. Ghem, K.M. O'Hara, T.A. Savard, and J.E. Thomas - Phys. Rev. A (1998)

S. Schneider and G. J. Milburn - Phys. Rev. A (1999)

T. Grotz, L. Heaney, and W. Strunz - Phys. Rev. A (2006)



Example 2

Application of intense laser pulse to the hydrogen-like atom

State space: $\mathfrak{h} = L^2(\mathbb{R}, \mathbb{C})$

Hamiltonian: $H = -\frac{1}{2} \frac{d^2}{dx^2} - \frac{1}{\sqrt{x^2 + a^2}} + xF(t)$, with

$$F(t) = F_0 \sin(\beta t + \delta) \cdot \begin{cases} \sin(\pi t / (2\tau)), & \text{if } t < \tau \\ 1, & \text{if } \tau \leq t \leq T - \tau \\ \cos^2(\pi(t + \tau - T) / (2\tau)), & \text{if } T - \tau \leq t \leq T \end{cases}.$$

$$L_1 = -i\eta x$$

$\beta, \eta, \delta \in \mathbb{R}$ and $a, F_0, \tau, T > 0$

$$X_t = X_0 + \int_0^t \left(-iH + \frac{1}{2} L_1^* L_1 \right) X_s ds + \int_0^t L_1 X_s dW_s^k$$

K.P. Singh and J.M. Rost - Phys. Rev. A (2007)



Example 3 (Quantum measurement)

Simultaneous measurement of position and momentum

State space: $\mathfrak{h} = L^2(\mathbb{R}, \mathbb{C})$

Hamiltonian: $H = -\alpha \frac{d^2}{dx^2} + \beta x^2$, with $\alpha \geq 0$ and $\beta \in \mathbb{R}$

$L_1 = \frac{\kappa}{\sigma} x$ and $L_2 = -i\kappa\sigma \frac{d}{dx}$, with $\kappa, \sigma \in]0, \infty[$

$$Y_t = Y_0 + \int_0^t G(Y_s) ds + \sum_{k=1}^2 \int_0^t L_k(Y_s) dB_s^k$$

- $G(y) = \left(-iH - \frac{1}{2} \sum_{k=1}^2 L_k^* L_k\right) y + \sum_{k=1}^2 (\operatorname{Re} \langle y, L_k y \rangle L_k y - \frac{1}{2} \operatorname{Re}^2 \langle y, L_k y \rangle y)$
- $L_k(y) = L_k y - \operatorname{Re} \langle y, L_k y \rangle y$

A.J. Scott and G.J. Milburn - Phys. Rev. A (2001)

J. Gough and A. Sobolev - Phys. Rev. A (2004)

A. Bassi and D. Dürr - Europhys. Lett. (2008)

Example 4

Quantum oscillator

State space: $\mathfrak{h} = l^2(\mathbb{Z}_+)$

$$H = i\beta_1 (a^\dagger - a) + \beta_2 N + \beta_3 (a^\dagger)^2 a^2$$

$$L_1 = \alpha_1 a, L_2 = \alpha_2 a^\dagger, L_3 = \alpha_3 N,$$

$$L_4 = \alpha_4 a^2, L_5 = \alpha_5 (a^\dagger)^2, L_6 = \alpha_6 N^2$$

$(e_n)_{n \in \mathbb{Z}_+}$: orthonormal basis of $l^2(\mathbb{Z}_+)$

$$a^\dagger e_n = \sqrt{n+1} e_{n+1}, \quad a e_n = \begin{cases} 0, & \text{if } n = 0 \\ \sqrt{n} e_{n-1}, & \text{if } n > 0 \end{cases}, \quad N = a^\dagger a$$

$$X_t = X_0 + \int_0^t \left(-iH + \frac{1}{2} \sum_{k=1}^6 L_k^* L_k \right) X_s ds + \sum_{k=1}^6 \int_0^t L_k X_s dW_s^k$$



Linear stochastic Schrödinger equation

Belavkin (Physics Letter A, 1989)

$$X_t(x) = x + \int_0^t G(s) X_s(x) ds + \sum_{k=1}^{\infty} \int_0^t L_k(s) X_s(x) dW_s^k \quad (1)$$

$$G(s) = -iH(s) - \frac{1}{2} \sum_{k=1}^{\infty} L_k(s)^* L_k(s)$$



Non-linear stochastic Schrödinger equation

Non-linear stochastic evolution equation

$$Y_t^y = y + \int_0^t G(s, Y_s^y) ds + \sum_{k=1}^{\infty} \int_0^t L_k(s, Y_s^y) dB_s^k \quad (2)$$

- $\|y\| = 1$
- B^1, B^2, \dots : independent Brownian motions.
- $L_k(s, x) = L_k(s)x - \operatorname{Re} \langle x, L_k(s)x \rangle x$.
- $G(s, x) = G(s)x + \sum_{k=1}^{\infty} \left(\operatorname{Re} \langle x, L_k(s)x \rangle L_k(s)x - \frac{1}{2} \operatorname{Re}^2 \langle x, L_k(s)x \rangle x \right)$.

Hypothesis 1

Hypothesis 1

Let C be a self-adjoint positive operator in \mathfrak{h} such that:

- For any $\xi \in L_C^2(\mathbb{P}, \mathfrak{h})$ and $T > 0$, the linear stochastic Schrödinger equation (1) has a unique strong C -solution on $[0, T]$ with initial datum ξ .
- For all $x \in \mathcal{D}(C)$ and $t \geq 0$,

$$2\Re \langle x, Gx \rangle + \sum_{k=1}^{\infty} \|L_k x\|^2 = 0.$$

- For any $x \in \mathcal{D}(C)$ and $t \geq 0$, $\|G(t)x\|^2 \leq K(t) \|x\|_C^2$.

$X_t(\xi)$ is strong C -solution iff

- $\mathbb{E} \|X_t(\xi)\|^2 \leq \mathbb{E} \|\xi\|^2$, $X_t(\xi) \in \mathcal{D}(C)$ a.s. and $\sup_{s \in [0, t]} \mathbb{E} \|CX_s(\xi)\|^2 < \infty$.
- $X_t(\xi) = \xi + \int_0^t G(s) \pi_C(X_s(\xi)) ds + \sum_{\ell=1}^{\infty} \int_0^t L_{\ell}(s) \pi_C(X_s(\xi)) dW_s^{\ell}$ \mathbb{P} -a.s.



Existence and uniqueness

Theorem (C.M. M and R. Rebolledo (Ann. Appl. Probab. (2008)) - F. Fagnola and C.M. M. (2011))

Let C satisfy Hypothesis 1.

Suppose that θ is a probability measure on $\mathfrak{B}(\mathfrak{h})$ such that $\theta(\mathcal{D}(C) \cap \{y \in \mathfrak{h} : \|y\| = 1\}) = 1$ and $\int_{\mathfrak{h}} \|Cy\|^2 \theta(dx) < \infty$.

Then the non-linear stochastic Schrödinger equation (2) has a unique C -solution $(\mathbb{Q}, (Y_t)_{t \geq 0}, (B_t)_{t \geq 0})$ with initial law θ .



Evolution of density operators

Quantum master equation

$$\begin{aligned}\frac{d}{dt}\rho_t(\varrho) &= \rho_t(\varrho) G(t)^* + G(t) \rho_t(\varrho) + \sum_{k=1}^{\infty} L_k(t) \rho_t(\varrho) L_k(t)^* \\ \rho_0(\varrho) &= \varrho\end{aligned}\quad (3)$$

Previous results

- Existence of minimal solution: Davies (Rep. Math. Phys., 1977)
- Uniqueness: Chebotarev, Fagnola (J. Funct. Anal., 1998)
- Regularity of solutions: Davies (Comm. math. phys., 1977, Neutron diffusion equation); Chebotarev, García and Quezada (Publ. Res. Inst. Math. Sci. Kokyuroku, 1998, General results); Arnold and Sparber (Commun. Math. Phys., 2004, Diffusion models with Hartree interaction)

- Is $\text{tr}(A\rho_t(\varrho))$ well-posed?



Regular density operator

Definition

Let C be an self-adjoint positive operator.

Then $\varrho \in \mathfrak{L}_{1,C}^+(\mathfrak{h})$ if and only if

- ϱ is a positive trace class operator in \mathfrak{h} .
- There is an orthonormal basis $(u_n)_{n \in \mathbb{Z}_+}$ of \mathfrak{h} and a sequence of non-negative real numbers $(\lambda_n)_{n \in \mathbb{Z}_+}$ such that:
 - $\varrho = \sum_{n \in \mathbb{Z}_+} \lambda_n |u_n\rangle \langle u_n|$.
 - $\sum_{n \in \mathbb{Z}_+} \lambda_n \|Cu_n\|^2 < +\infty$

Characterizations of $\mathfrak{L}_{1,C}^+(\mathfrak{h})$: Chebotarev, Garcia and Quezada (1998)

Lemma

ϱ is C -regular if and only if there exists $\xi \in L_C^2(\mathbb{P}, \mathfrak{h})$ for which $\varrho = \mathbb{E} |\xi\rangle \langle \xi|$.



Evolution of density operators

Theorem

Suppose that Hypothesis 1 holds. Then, for every $t \geq 0$ there exists a unique operator ρ_t belonging to $\mathfrak{L}(\mathfrak{L}_1(\mathfrak{h}))$ such that for each C -regular operator ϱ we have

$$\rho_t(\varrho) = \mathbb{E} \left| Y_t^\xi \right\rangle \left\langle Y_t^\xi \right| = \mathbb{E} \left| X_t^\xi \right\rangle \left\langle X_t^\xi \right|,$$

whenever ξ is an arbitrary random variable satisfying $\xi \in L_C^2(\mathbb{P}, \mathfrak{h})$ and $\varrho = \mathbb{E} |\xi\rangle \langle \xi|$.

Theorem

Under Hypothesis 1,

$$\rho_t(\mathfrak{L}_1^+(\mathfrak{h})) \subset \mathfrak{L}_1^+(\mathfrak{h}) \text{ and } \rho_t(\mathfrak{L}_{1,C}^+(\mathfrak{h})) \subset \mathfrak{L}_{1,C}^+(\mathfrak{h}).$$

Mean values

Lemma

Suppose that:

- $\varrho = \mathbb{E} |\xi\rangle\langle\xi|$ for $\xi \in L^2_C(\mathbb{P}, \mathfrak{h})$.
- $A \in \mathfrak{L}((\text{Dom}(C), \langle \cdot, \cdot \rangle_C), \mathfrak{h})$.

Then $A\varrho = \mathbb{E} |A\xi\rangle\langle\xi|$ and $\text{tr}(A\varrho) = \mathbb{E} \langle\xi, A\xi\rangle$.

If in addition $\text{Dom}(C) \subset \text{Dom}(A^*)$, we have

- $\varrho A = \mathbb{E} |\xi\rangle\langle A^*\xi|$.
- $\text{tr}(\varrho A) = \mathbb{E} \langle\xi, A\xi\rangle$.



Existence and uniqueness of solutions for QMEs

Theorem (C.M. Ann. Probab. (to appear))

Consider the autonomous case. Let Hypothesis 1 hold. Then for any $A \in \mathfrak{L}(\mathfrak{h})$ and $t \geq 0$,

$$\frac{d}{dt} \operatorname{tr}(A \rho_t(\varrho)) = \operatorname{tr} \left(A \left(G \rho_t(\varrho) + \rho_t(\varrho) G^* + \sum_{k=1}^{\infty} L_k \rho_t(\varrho) L_k^* \right) \right). \quad (4)$$

Moreover, $(\rho_t)_{t \geq 0}$ is the unique semigroup of bounded operators on $\mathfrak{L}_1(\mathfrak{h})$ such that:

- i) $\sup_{t \in [0, T]} \|\rho_t\|_{\mathfrak{L}(\mathfrak{L}_1(\mathfrak{h}))} < \infty$.
- ii) For each $x \in \operatorname{Dom}(C)$, the function $t \mapsto \operatorname{tr}(\rho_t(|x\rangle\langle x|)A)$ is continuous provided $A \in \mathfrak{L}(\mathfrak{h})$.
- iii) Relation (4) holds with $\varrho = |x\rangle\langle x|$ whenever $x \in \operatorname{Dom}(C)$.



Hypothesis 2 (non-explosion condition)

Let C be a self-adjoint positive operator in \mathfrak{h} with the properties:

- For any $x \in \mathcal{D}(C)$, $\|G(t)x\|^2 \leq K(t)\|x\|_C^2$.
- For all $x \in \mathcal{D}(C)$, $\|L_k(t)x\|^2 \leq K(t)\|x\|_C^2$.
- There exists $\alpha \geq 0$ and a core \mathfrak{D}_1 of C^2 such that for all $x \in \mathfrak{D}_1$,

$$2\Re \langle C^2 x, G(t)x \rangle + \sum_{k=1}^{\infty} \|CL_k(t)x\|^2 \leq \alpha(t)\|x\|_C^2.$$

- There exist a core \mathfrak{D}_2 of C such that for any x in \mathfrak{D}_2 ,

$$2\Re \langle x, G(t)x \rangle + \sum_{k=1}^{\infty} \|L_k(t)x\|^2 \leq 0.$$



Theorem (C.M. M., F. Fagnola, 2011)

Assume that Hypothesis 2 holds.

Let $\xi \in L_C^2(\mathbb{P}, \mathfrak{h})$.

Then, the linear stochastic Schrödinger equation (1) has a unique strong C -solution $(X_t(\xi))_{t \geq 0}$ with initial datum ξ .

Moreover,

$$\mathbb{E} \|CX_t(\xi)\|^2 \leq \exp(\alpha t) \left(\mathbb{E} \|C\xi\|^2 + \alpha t \mathbb{E} \|\xi\|^2 + \beta t \right).$$

Model 1

Consider $\mathfrak{h} = L^2(\mathbb{R}^d, \mathbb{C})$. Let the Hamiltonian be

$$H(t) = -\alpha\Delta + i \sum_{j=1}^d \left(A^j(t, \cdot) \partial_j + \partial_j A^j(t, \cdot) \right) + V(t, \cdot),$$

where $t \geq 0$, $\alpha \geq 0$, and V, A^1, \dots, A^d are real-valued measurable smooth functions on $[0, +\infty[\times \mathbb{R}^d$.

For a given $m \in \mathbb{N}$ and for all $t \geq 0$ choose

$$L_\ell(t) = \begin{cases} \sum_{k=1}^d \sigma_{\ell k}(t, \cdot) \partial_k + \eta_\ell(t, \cdot), & \text{if } 1 \leq \ell \leq m \\ 0, & \text{if } \ell > m \end{cases},$$

where $\sigma_{\ell k}, \eta_\ell : [0, +\infty[\times \mathbb{R}^d \rightarrow \mathbb{C}$ are complex-valued measurable smooth functions.



Hypothesis 3

Adopt Model 1. Define $G(t) = -iH(t) - \frac{1}{2} \sum_{\ell=1}^m L_{\ell}^*(t) L_{\ell}(t)$.

(H3.1)

Suppose that: $V(t, \cdot) \in C^2(\mathbb{R}^d, \mathbb{R})$, $A^j(t, \cdot) \in C^3(\mathbb{R}^d, \mathbb{R})$,

$$\max \{ |V(t, x)|, |\Delta V(t, x)|, |\partial_j(\Delta A^j)| \} \leq K(t) (1 + |x|^2),$$

$$\max \{ |\partial_j V(t, x)|, |A^j(t, x)|, |(\partial_{j'} \partial_j A^j)(t, x)| \} \leq K(t) (1 + |x|)$$

$$|\partial_{j'} A^j(t, x)| \leq K(t)$$

$$H(t) = -\alpha \Delta + i \sum_{j=1}^d \left(A^j(t, \cdot) \partial_j + \partial_j A^j(t, \cdot) \right) + V(t, \cdot)$$



Hypothesis 3

(H3.2)

$|\sigma_{\ell k}(t, \cdot)| \leq K(t)$,
 $\eta_{\ell}(t, \cdot) \in C^3(\mathbb{R}^d, \mathbb{C})$ and the absolute values of all the partial derivatives of $\eta_{\ell}(t, \cdot)$ from the first up to the third order are bounded by $K(t)$.

At least one of the following conditions holds:

- $|\eta_{\ell}(t, \cdot)| \leq K(t)$, $\sigma_{\ell k}(t, \cdot) \in C^3(\mathbb{R}^d, \mathbb{C})$, and the absolute values of all partial derivatives of $\sigma_{\ell k}(t, \cdot)$ up to the third order are dominated by $K(t)$.
- $(t, x) \mapsto \sigma_{\ell k}(t, x)$ does not depend on x and $|\eta_{\ell}(t, 0)| \leq K(t)$.

$$L_{\ell}(t) = \sum_{k=1}^d \sigma_{\ell k}(t, \cdot) \partial_k + \eta_{\ell}(t, \cdot)$$



Theorem (CMM, F. Fagnola (2011))

Suppose that Hypothesis 3 holds.

Set $C = -\Delta + |x|^2$.

Let ξ be a \mathfrak{F}_0 -measurable random variable taking values in $L^2(\mathbb{R}^d, \mathbb{C})$ such that $\mathbb{E} \|\xi\|^2 = 1$ and $\mathbb{E} \|C\xi\|^2 < \infty$.

Then the linear SSE (1) has a unique strong C -solution with initial datum ξ . Moreover, $\mathbb{E} \|X_t(\xi)\|^2 = \|\xi\|^2$ for all $t > 0$.

Hypothesis 4

Let C satisfy Hypotheses 1. Suppose that:

For all $t \geq 0$ and any x belonging to a core of C ,

$$\sum_{\ell=1}^{\infty} \|C^{1/2} L_{\ell}(t) x\|^2 \leq K(t) \|x\|_C^2.$$

Let $A = B_1^* B_2$, where B_1, B_2 are operators in \mathfrak{h} such that:

- For all $x \in \mathcal{D}(C^{1/2})$, $\max\{\|B_1 x\|^2, \|B_2 x\|^2\} \leq K \|x\|_{C^{1/2}}^2$.
- $\max\{\|Ax\|^2, \|A^* x\|^2\} \leq K \|x\|_C^2$ whenever $x \in \mathcal{D}(C)$.



Theorem F. Fagnola and C.M. M. (2012)

Let Hypothesis 4 hold.

Assume the existence and uniqueness of a strong C -solution to the linear SSE (1) with initial datum $\xi \in L^2_C(\mathbb{P}; \mathfrak{h})$ on any bounded interval.

Then, for all $t \geq 0$ we have

$$\begin{aligned} \mathbb{E} \langle X_t(\xi), AX_t(\xi) \rangle &= \mathbb{E} \langle \xi, A\xi \rangle + \int_0^t \mathbb{E} \langle A^* X_s(\xi), GX_s(\xi) \rangle ds \\ &\quad + \int_0^t \mathbb{E} \langle GX_s(\xi), AX_s(\xi) \rangle ds \\ &\quad + \int_0^t \left(\sum_{\ell=1}^{\infty} \mathbb{E} \langle B_1 L_\ell X_s(\xi), B_2 L_\ell X_s(\xi) \rangle \right) ds \end{aligned}$$



Theorem F. Fagnola and C.M. M. (2012)

Assume the context of Model 1, together with Hypothesis 3.

For any $j = 1, 2$, let B_j be either $\partial_k [a_j]$, $[b_j] \partial_k$ or $[c_j]$, where $k = 1, \dots, d$, $a_j \in C^2(\mathbb{R}^d, \mathbb{R})$ and $b_j, c_j \in C^1(\mathbb{R}^d, \mathbb{R})$.

Suppose that:

$$\max \left\{ |a_j(x)|, |b_j(x)| \right\} \leq K,$$

$$\max \left\{ |c_j(x)|, |\partial_l a_j(x)|, |\partial_l b_j(x)| \right\} \leq K(1 + |x|), \text{ and}$$

$$\max \left\{ |\partial_l c_j(x)|, |\partial_k \partial_l a_j(x)| \right\} \leq K(1 + |x|^2).$$

If $A = B_1^* B_2$ and $\xi \in L^2_{-\Delta + |x|^2}(\mathbb{P}; \mathfrak{h})$, then

$$\begin{aligned} \mathbb{E} \langle X_t(\xi), AX_t(\xi) \rangle &= \mathbb{E} \langle \xi, A\xi \rangle + \int_0^t \mathbb{E} \langle A^* X_s(\xi), GX_s(\xi) \rangle ds \\ &\quad + \int_0^t \mathbb{E} \langle GX_s(\xi), AX_s(\xi) \rangle ds \\ &\quad + \int_0^t \left(\sum_{\ell=1}^{\infty} \mathbb{E} \langle B_1 L_\ell X_s(\xi), B_2 L_\ell X_s(\xi) \rangle \right) ds \end{aligned}$$

Example 1

State space: $\mathfrak{h} = L^2(\mathbb{R}, \mathbb{C})$

Hamiltonian: $H = -\frac{1}{2M} \frac{d^2}{dx^2} + \frac{1}{2} M \omega^2 x^2$, with $M > 0$ and $\omega \in \mathbb{R}$

$L_1 = -i\eta x$, with $\eta > 0$.

Theorem F. Fagnola and C.M. M. (2012)

In Example 1, for all $t \geq 0$ we have

$$\mathbb{E} \langle X_t(\xi), H X_t(\xi) \rangle = \mathbb{E} \langle \xi, H \xi \rangle + \frac{1}{2M} \eta^2 t.$$

Hypothesis 5

Hypothesis 5

Suppose that there exist: (i) a self-adjoint positive operator D in \mathfrak{h} satisfying Hypothesis 1; and (ii) a probability measure Γ on $\mathfrak{B}(\mathfrak{h})$ such that:

- For any $t \geq 0$ and $A \in \mathfrak{B}(\mathfrak{h})$, $\Gamma(A) = \int_{\mathfrak{h}} P_t(x, A) \Gamma(dx)$
- $\Gamma(\text{Dom}(D) \cap \{x \in \mathfrak{h} : \|x\| = 1\}) = 1$
- $\int_{\mathfrak{h}} \|Dz\|^2 \Gamma(dz) < \infty$

Here $P_t(x, A) = \begin{cases} \mathbb{Q}_x(Y_t^x \in A), & x \in \text{Dom}(D) \\ \delta_x(A), & x \notin \text{Dom}(D) \end{cases}$

Sufficient condition: AAP (2008) C.M.M - R. Rebolledo



Regular invariant density operator

Theorem

Let D satisfy Hypothesis 5.

Then, there exists a D -regular operator ϱ_∞ such that

$$\rho_t(\varrho_\infty) = \varrho_\infty$$

for all $t \geq 0$.



Example 4 (quantum oscillator)

State space: $\mathfrak{h} = l^2(\mathbb{Z}_+)$

$(e_n)_{n \in \mathbb{Z}_+}$: orthonormal basis of $l^2(\mathbb{Z}_+)$

$$a^\dagger e_n = \sqrt{n+1} e_{n+1}, \quad ae_n = \begin{cases} 0, & \text{if } n = 0 \\ \sqrt{n} e_{n-1}, & \text{if } n > 0 \end{cases}$$

$$N = a^\dagger a$$

$$H = i\beta_1 (a^\dagger - a) + \beta_2 N + \beta_3 (a^\dagger)^2 a^2$$

$$L_1 = \alpha_1 a, L_2 = \alpha_2 a^\dagger, L_3 = \alpha_3 N,$$

$$L_4 = \alpha_4 a^2, L_5 = \alpha_5 (a^\dagger)^2, L_6 = \alpha_6 N^2$$

Unbounded observables

- N : number of photons
- $i(a^\dagger - a)/\sqrt{2}$: The position operator
- $(a^\dagger + a)/\sqrt{2}$: The momentum operator

Example 4

$$H = i\beta_1 (a^\dagger - a) + \beta_2 N + \beta_3 (a^\dagger)^2 a^2, \quad L_1 = \alpha_1 a, \quad L_2 = \alpha_2 a^\dagger, \\ L_3 = \alpha_3 N, \quad L_4 = \alpha_4 a^2, \quad L_5 = \alpha_5 (a^\dagger)^2, \quad L_6 = \alpha_6 N^2$$

Theorem

In the set-up of Example 4 we assume

$$|\alpha_4| \geq |\alpha_5|.$$

If $p \geq 4$, then there exists a unique N^p -regular solution to the quantum master equation (3).

In addition, there exists a N^p -regular operator ϱ_∞ which is invariant for (3) provided that either

$$|\alpha_4| > |\alpha_5| \text{ or}$$

$$|\alpha_4| = |\alpha_5| \text{ with } |\alpha_2|^2 - |\alpha_1|^2 + 4(2p+1)|\alpha_4|^2 < 0.$$



Conclusion

- We can study theoretical properties of the quantum Markovian master equations with the help of stochastic Schrödinger equations.
- We obtain an Ehrenfest's-type theorem for open quantum systems.
- We can prove rigorously the heating of ion traps in a simple model.
- In many physical situations, there exists a regular stationary solution for the quantum master equations.



Thank you very much!

